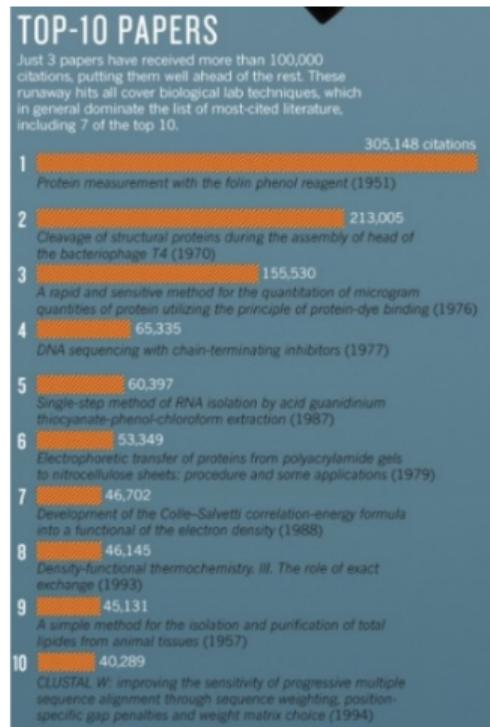


# Density Functional Theory



## Density Functional Theory extremely influential

- 1926: Schrödinger's equation
- 1927: Thomas-Fermi
- 1965: Kohn-Sham

Known: Obtaining observables from quantum mechanics (Born-Oppenheimer)

$$\hat{H} = \hat{H} (Z_i, \mathbf{R}_i, N_e, \sigma) \quad (50)$$

Number of electrons      Spin  
Nuclear charges      Nuclear coordinates

### Steps:

- System defines Hamiltonian  $\hat{H}$
- Solving Schrödinger's equation yields wave function  $\hat{H} \Psi = E \Psi$
- Operator yields observable  $\langle \hat{A} \rangle = \langle \Psi | \hat{A} | \Psi \rangle$

$$\hat{H} = \hat{T} + \hat{V} = \hat{T}_e + \hat{V}_{en} + \hat{V}_{ee} + \hat{V}_{nn}$$

Electron-electron interaction      Nuclei-nuclei  
 Kinetic energy      Potential energy      Electron-nuclei interaction  
 Kinetic energy electrons

(51)

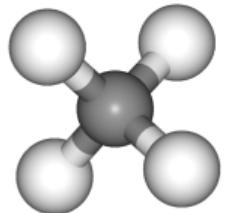
with

$$\hat{H} = - \sum_i \frac{\hbar^2}{2m_e} \nabla_{r_i}^2 - \sum_{i,j} \frac{Z_i e^2}{4\pi\epsilon_0 |\mathbf{R}_i - \mathbf{r}_j|} + \sum_{i,j>i} \frac{e^2}{4\pi\epsilon_0 |\mathbf{r}_i - \mathbf{r}_j|} + \sum_{i,j>i} \frac{Z_i Z_j e^2}{4\pi\epsilon_0 |\mathbf{R}_i - \mathbf{R}_j|}$$
(52)

Resulting wave function

$$\Psi = \Psi(\mathbf{r}_1, \dots, \mathbf{r}_n)$$

has  $3N_e$  dimensions!



## Example

- Methane, i.e. 10 electrons
- Grid with 5 points along each dimension

## Wave function

- $5^{3 \cdot 10}$  grid points
- 6 ZB
- 7 times ALL data

Replace  $3N$ -dimensional wave function  $\Psi$  with a 3-dimensional electron density  $\rho$

$$\Psi = \Psi(\mathbf{r}_1, \dots, \mathbf{r}_n) \longrightarrow \rho = \rho(\mathbf{r}) \quad (53)$$

E.B. Wilson<sup>1</sup>: this contains the same information

$$\rho \rightarrow Z_i, \mathbf{R}_i, N_e, \sigma \rightarrow \hat{H} \rightarrow \Psi$$

$$\begin{array}{c} \text{Number of electrons} \quad \text{Spin} \\ \hat{H}(Z_i, \mathbf{R}_i, N_e, \sigma) \rightarrow \Psi \\ \text{Nuclear charges} \quad \text{Nuclear coordinates} \end{array} \quad (54)$$

## Tricks:

- Coordinates: local maxima in  $\rho$
- Nuclear charges: Kato's Theorem<sup>2</sup>
- Number electrons and spin:  $N_e = \int d\mathbf{r} \rho(\mathbf{r})$

*The ground state electron density uniquely defines the external potential.<sup>1</sup>*

Proof through contradiction: Let's assume, one density would be valid for two external potentials  $V_i$

$$E_1 = \langle \Psi_1 | \hat{H}_1 | \Psi_1 \rangle = \int V_1(\mathbf{r}) \rho(\mathbf{r}) d\mathbf{r} + \langle \Psi_1 | \hat{T} + \hat{U} | \Psi_1 \rangle \quad (55)$$

Annotations for the equation:

- External potential: Points to  $V_1(\mathbf{r})$
- Electron density for both systems: Points to  $\rho(\mathbf{r})$
- Wave function for potential 1: Points to  $\Psi_1$
- Energy for potential 1: Points to  $E_1$
- Universal (HK) functional: Points to  $\hat{T} + \hat{U}$
- Kinetic energy: Points to  $\hat{T}$

Variational principle  $\Rightarrow$  For all other wave functions, energy is strictly higher

$$E_1 < \langle \Psi_2 | \hat{H}_1 | \Psi_2 \rangle = \langle \Psi_2 | \hat{H}_2 | \Psi_2 \rangle + \langle \Psi_2 | \hat{H}_1 - \hat{H}_2 | \Psi_2 \rangle = E_2 + \int [V_1(\mathbf{r}) - V_2(\mathbf{r})] \rho(\mathbf{r}) d\mathbf{r} \quad (56)$$

Hamiltonian system 2  
Hamiltonian system 1  
Wave function for potential 2

Due to symmetry between indices 1 and 2:

$$E_2 < E_1 + \int [V_2(\mathbf{r}) - V_1(\mathbf{r})] \rho(\mathbf{r}) d\mathbf{r} \quad (57)$$

Contradiction from summation

$$E_1 + E_2 < E_1 + E_2 \quad (58)$$

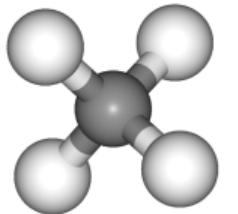
*The ground state electron density has the lowest energy of all densities.<sup>1</sup>*

Proof: For other wave functions, the energy is strictly larger.

$$E_0 = \langle \Psi_0 | \hat{H} | \Psi_0 \rangle, \quad E_0 < \langle \Psi | \hat{H} | \Psi \rangle \quad (\Psi \neq \Psi_0) \quad (59)$$

Since each wave function has exactly one energy.

$$E_0 = E[\rho_0] < E[\rho], \quad \rho \neq \rho_0 \quad (60)$$



## Example

- Methane, i.e. 10 electrons
- Grid with 5 points along each dimension

## Wave function

- $5^{3 \cdot 10}$  grid points
- 6 ZB
- 7 times ALL data

## Electron density

- $5^3$  grid points
- 1 KB
- half a page of text

## What we have

- Unknown universal functional  $F[\rho]$  (universal: does not depend on geometry)
- Trial densities  $\rho$

## What we need

- some functional  $F$
- procedure to find optimal  $\rho_0$

## System

- Take geometry, obtain  $V_{\text{ext}}$
- Choose functional and basis set (see later)
- Minimize  $E[\rho]$  by varying density self-consistently

$$E[\rho] = F[\rho] + \int V_{\text{ext}} \rho d\mathbf{r} \quad (61)$$

Challenge: Hard to find  $F$ . Kohn and Sham proposed<sup>1</sup>

$$F[\rho] = T_s[\rho] + E_H[\rho] + E_{\text{xc}}[\rho] \quad (62)$$

Hartree energy functional      Exchange correlation functional  
 ↓                                    ↓  
 Kinetic energy non-interacting particles

$$T_s[\rho] = \langle \Phi[\rho] | \hat{T} | \Phi[\rho] \rangle; \quad E_H[\rho] = \frac{1}{2} \iint \frac{\rho(\mathbf{r}_1) \rho(\mathbf{r}_2)}{|\mathbf{r}_1 - \mathbf{r}_2|} d\mathbf{r}_1 d\mathbf{r}_2 \quad (63)$$

Kohn-Sham wave function      ↓  
 ↓                                    ↓  
 Kinetic energy operator  $-\nabla^2/2$

$$E[\{\varphi_i\}] = \sum_{i=1}^N \int \varphi_i^*(\mathbf{r}) \left( -\frac{1}{2} \nabla^2 + V_{\text{ext}}(\mathbf{r}) \right) \varphi_i(\mathbf{r}) d\mathbf{r} + E_{\mathbb{H}}[\rho] + E_{\text{xc}}[\rho] \quad (64)$$

↑ Kohn-Sham orbitals

Where the orbitals are found by solving (self-consistently)

$$\left( -\frac{1}{2} \nabla^2 + V_{\text{ext}}(\mathbf{r}) + \frac{\delta [E_{\mathbb{H}}[\rho] + E_{\text{xc}}[\rho]]}{\delta \rho(\mathbf{r})}(\mathbf{r}) \right) \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r}) \quad (65)$$

↑ Kohn-Sham orbital energies

$$\frac{v_{\mathbb{H}}(\mathbf{r})}{\text{↑ Hartree potential}} \equiv \frac{\delta E_{\mathbb{H}}[\rho]}{\delta \rho(\mathbf{r})}(\mathbf{r}) = \int \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}'; \frac{v_{\text{xc}}(\mathbf{r})}{\text{↑ Exchange-correlation potential}} \equiv \frac{\delta E_{\text{xc}}[\rho]}{\delta \rho(\mathbf{r})}(\mathbf{r}) \quad (66)$$

How to solve? See Hartree-Fock (later)

# Recent developments in LIBXC – A comprehensive library of functionals for density functional theory



Susi Lehtola <sup>a,1</sup>, Conrad Steigemann <sup>b</sup>, Micael J.T. Oliveira <sup>c,\*</sup>, Miguel A.L. Marques <sup>b</sup>

<sup>a</sup> Chemical Sciences Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, United States

<sup>b</sup> Institut für Physik, Martin-Luther-Universität Halle-Wittenberg, D-06099 Halle, Germany

<sup>c</sup> Max Planck Institute for the Structure and Dynamics of Matter, Luruper Chaussee 149, 22761 Hamburg, Germany

---

## ARTICLE INFO

---

### Article history:

Received 30 June 2017

Received in revised form 6 November 2017

Accepted 7 November 2017

---

### Keywords:

Density functional theory

Exchange–correlation

Local density approximations

Generalized gradient approximations

meta-GGA approximations

---

---

## ABSTRACT

LIBXC is a library of exchange–correlation functionals for density-functional theory. We are concerned with semi-local functionals (or the semi-local part of hybrid functionals), namely local-density approximations, generalized-gradient approximations, and meta-generalized-gradient approximations. Currently we include around 400 functionals for the exchange, correlation, and the kinetic energy, spanning more than 50 years of research. Moreover, LIBXC is by now used by more than 20 codes, not only from the atomic, molecular, and solid-state physics, but also from the quantum chemistry communities.

© 2017 The Authors. Published by Elsevier B.V. This is an open access article under the CC BY license (<http://creativecommons.org/licenses/by/4.0/>).

## Idea

- Assumes (almost) same density everywhere
- Exact for the uniform electron gas *Jellium*
- Used for materials, deprecated
- Strictly local

$$E_{xc}^{\text{LDA}} = \int d\mathbf{r} \rho(\mathbf{r}) \epsilon_{xc}^{\text{LDA}}[\rho(\mathbf{r})] \quad (67)$$

Popular choice: VWN

## Idea

- Includes local gradient
- *semi-local* nature: gradient tells a bit about close proximity
- Widely used, both for materials and molecules

$$E_{xc}^{\text{GGA}} = \int d\mathbf{r} \rho(\mathbf{r}) \epsilon_{xc}^{\text{GGA}}(\rho(\mathbf{r}), \nabla \rho(\mathbf{r})) \quad (68)$$

Popular choice: PBE

### Idea

- Next derivative and/or kinetic energy density  $\tau$
- Sensitive to numerical integration through Laplacian
- Default settings sometimes not good enough

$$E_{xc}^{\text{GGA}} = \int d\mathbf{r} \rho(\mathbf{r}) \epsilon_{xc}^{\text{GGA}}(\rho(\mathbf{r}), \nabla \rho(\mathbf{r}), \nabla^2 \rho(\mathbf{r}), \tau), \quad \tau \equiv \sum_i^{\text{occ}} |\nabla \psi_i(\mathbf{r})|^2 \quad (69)$$

Popular choices: SCAN (built on constraints!), M06, TPSS

**Idea**

- Commonly improved band gaps and relative energies between spin states
- Popular choices: B3LYP, PBE0

$$E_{xc}^{\text{hyb}} = a E_x^{\text{HF}} + (1 - a) E_x^{\text{DFT}} + E_c^{\text{DFT}} \quad (70)$$

$$E_x^{\text{HF}} = -\frac{1}{2} \sum_{i,j} \iint \psi_i^*(\mathbf{r}_1) \psi_j^*(\mathbf{r}_2) \frac{1}{r_{12}} \psi_j(\mathbf{r}_1) \psi_i(\mathbf{r}_2) \, d\mathbf{r}_1 \, d\mathbf{r}_2 \quad (71)$$

**Idea**

- Computationally more efficient
- Popular choice: HSE06

$$\frac{1}{r_{12}} = \frac{1 - [\alpha + \beta(1 - \omega_{\text{RSF}}(\gamma, r_{12}))]}{r_{12}} + \frac{\alpha + \beta(1 - \omega_{\text{RSF}}(\gamma, r_{12}))}{r_{12}} \quad (72)$$

Can be tuned in many cases to specific system based on Koopman's condition

$$-\varepsilon_{\text{HOMO}}(N_e) = \text{IE}(N_e) = E(N_e - 1) - E(N_e) \quad (73)$$

## Implications

- Replace the  $3N$ -dimensional wave function with 3D electron density.
- There is one functional (function of functions) yielding the energy, hence the name density functional theory:  $E[\rho]$
- This functional is...
  - Universal: Does not depend on the system!
  - Unknown: Topic of current research
- We can obtain the electron density by minimizing the energy.