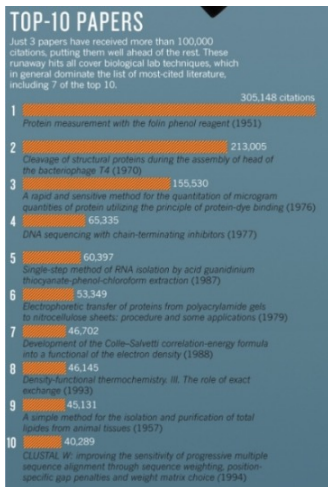


Density Functional Theory



Density Functional Theory extremely influential

- 1926: Schrödinger's equation
- 1927: Thomas-Fermi
- 1965: Kohn-Sham

Known: Obtaining observables from quantum mechanics (Born-Oppenheimer)

$$\hat{H} = \hat{H}(\underbrace{Z_i}_{\text{Nuclear charges}}, \underbrace{\mathbf{R}_i}_{\text{Nuclear coordinates}}, \underbrace{N_e}_{\text{Number of electrons}}, \underbrace{\sigma}_{\text{Spin}}) \quad (50)$$

Steps:

- System defines Hamiltonian \hat{H}
- Solving Schrödinger's equation yields wave function $\hat{H} \Psi = E \Psi$
- Operator yields observable $\langle \hat{A} \rangle = \langle \Psi | \hat{A} | \Psi \rangle$

$$\hat{H} = \hat{T} + \hat{V} = \hat{T}_e + \hat{V}_{en} + \hat{V}_{ee} + \hat{V}_{nn} \quad (51)$$

Electron-electron interaction
 Nuclei-nuclei
 Kinetic energy
 Potential energy
 Electron-nuclei interaction
 Kinetic energy electrons

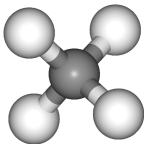
with

$$\hat{H} = - \sum_i \frac{\hbar^2}{2m_e} \nabla_{r_i}^2 - \sum_{i,j} \frac{Z_i e^2}{4\pi\epsilon_0 |\mathbf{R}_i - \mathbf{r}_j|} + \sum_{i,j>i} \frac{e^2}{4\pi\epsilon_0 |\mathbf{r}_i - \mathbf{r}_j|} + \sum_{i,j>i} \frac{Z_i Z_j e^2}{4\pi\epsilon_0 |\mathbf{R}_i - \mathbf{R}_j|} \quad (52)$$

Resulting wave function

$$\Psi = \Psi(\mathbf{r}_1, \dots, \mathbf{r}_n)$$

has $3N_e$ dimensions!



Example

- Methane, i.e. 10 electrons
- Grid with 5 points along each dimension

Wave function

- $5^{3 \cdot 10}$ grid points
- 6 ZB
- 7 times ALL data

Replace 3N-dimensional wave function Ψ with a 3-dimensional electron density ρ

$$\Psi = \Psi(\mathbf{r}_1, \dots, \mathbf{r}_n) \longrightarrow \rho = \rho(\mathbf{r}) \quad (53)$$

E.B. Wilson¹: this contains the same information

$$\rho \rightarrow Z_i, \mathbf{R}_i, N_e, \sigma \rightarrow \hat{H} \rightarrow \Psi$$

$$\begin{array}{c} \text{Number of electrons} \quad \text{Spin} \\ \hat{H}(Z_i, \mathbf{R}_i, N_e, \sigma) \rightarrow \Psi \\ \text{Nuclear charges} \quad \text{Nuclear coordinates} \end{array} \quad (54)$$

Tricks:

- Coordinates: local maxima in ρ
- Nuclear charges: Kato's Theorem²
- Number electrons and spin: $N_e = \int d\mathbf{r} \rho(\mathbf{r})$

¹ | P.O. Löwdin, Int. J. Quant. Chem. 1966. ² | T. Kato, Commun. Pure Appl. Math., 1957.

*The ground state electron density uniquely defines the external potential.*¹

Proof through contradiction: Let's assume, one density would be valid for two external potentials V_i

$$E_1 = \langle \Psi_1 | \hat{H}_1 | \Psi_1 \rangle = \int V_1(\mathbf{r}) \rho(\mathbf{r}) d\mathbf{r} + \langle \Psi_1 | \hat{T} + \hat{U} | \Psi_1 \rangle \quad (55)$$

External potential $V_1(\mathbf{r})$ Electron density for both systems $\rho(\mathbf{r})$
 Wave function for potential 1 Ψ_1 Universal (HK) functional \hat{U}
 Energy for potential 1 E_1 Kinetic energy \hat{T}

¹ | P. Hohenberg, W. Kohn, Phys. Rev. 1964.

Variational principle \Rightarrow For all other wave functions, energy is strictly higher

$$E_1 < \langle \Psi_2 | \hat{H}_1 | \Psi_2 \rangle = \langle \Psi_2 | \hat{H}_2 | \Psi_2 \rangle + \langle \Psi_2 | \hat{H}_1 - \hat{H}_2 | \Psi_2 \rangle = E_2 + \int [V_1(\mathbf{r}) - V_2(\mathbf{r})] \rho(\mathbf{r}) d\mathbf{r} \quad (56)$$

Hamiltonian system 2
Hamiltonian system 1

Wave function for potential 2

Due to symmetry between indices 1 and 2:

$$E_2 < E_1 + \int [V_2(\mathbf{r}) - V_1(\mathbf{r})] \rho(\mathbf{r}) d\mathbf{r} \quad (57)$$

Contradiction from summation

$$E_1 + E_2 < E_1 + E_2 \quad (58)$$

*The ground state electron density has the lowest energy of all densities.*¹

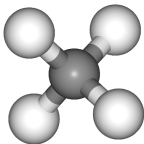
Proof: For other wave functions, the energy is strictly larger.

$$E_0 = \langle \Psi_0 | \hat{H} | \Psi_0 \rangle, \quad E_0 < \langle \Psi | \hat{H} | \Psi \rangle \quad (\Psi \neq \Psi_0) \quad (59)$$

Since each wave function has exactly one energy.

$$E_0 = E[\rho_0] < E[\rho], \quad \rho \neq \rho_0 \quad (60)$$

¹ | P. Hohenberg, W. Kohn, Phys. Rev. 1964.



Example

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- Grid with 5 points along each dimension

Wave function

- $5^{3 \cdot 10}$ grid points
- 6 ZB
- 7 times ALL data

Electron density

- 5^3 grid points
- 1 KB
- half a page of text

What we have

- Unknown universal functional $F[\rho]$ (universal: does not depend on geometry)
- Trial densities ρ

What we need

- some functional F
- procedure to find optimal ρ_0

System

- Take geometry, obtain V_{ext}
- Choose functional and basis set (see later)
- Minimize $E[\rho]$ by varying density self-consistently

$$E[\rho] = F[\rho] + \int V_{\text{ext}} \rho d\mathbf{r} \quad (61)$$

Challenge: Hard to find F . Kohn and Sham proposed¹

Hartree energy functional Exchange correlation functional

$$F[\rho] = T_s[\rho] + E_H[\rho] + E_{\text{xc}}[\rho] \quad (62)$$

Kinetic energy non-interacting particles

Kohn-Sham wave function

$$T_s[\rho] = \langle \Phi[\rho] | \hat{T} | \Phi[\rho] \rangle; \quad E_H[\rho] = \frac{1}{2} \iint \frac{\rho(\mathbf{r}_1) \rho(\mathbf{r}_2)}{|\mathbf{r}_1 - \mathbf{r}_2|} d\mathbf{r}_1 d\mathbf{r}_2 \quad (63)$$

Kinetic energy operator $-\nabla^2/2$

¹ | W. Kohn and L. J. Sham, Phys. Rev., 1965.

$$E[\{\varphi_i\}] = \sum_{i=1}^N \int \varphi_i^*(\mathbf{r}) \left(-\frac{1}{2} \nabla^2 + V_{\text{ext}}(\mathbf{r}) \right) \varphi_i(\mathbf{r}) d\mathbf{r} + E_{\text{H}}[\rho] + E_{\text{xc}}[\rho] \quad (64)$$

↑ Kohn-Sham orbitals

Where the orbitals are found by solving (self-consistently)

$$\left(-\frac{1}{2} \nabla^2 + V_{\text{ext}}(\mathbf{r}) + \frac{\delta [E_{\text{H}}[\rho] + E_{\text{xc}}[\rho]]}{\delta \rho(\mathbf{r})}(\mathbf{r}) \right) \varphi_i(\mathbf{r}) = \epsilon_i \varphi_i(\mathbf{r}) \quad (65)$$

↑ Kohn-Sham orbital energies

$$v_{\text{H}}(\mathbf{r}) \equiv \frac{\delta E_{\text{H}}[\rho]}{\delta \rho(\mathbf{r})}(\mathbf{r}) = \int \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}'; \quad v_{\text{xc}}(\mathbf{r}) \equiv \frac{\delta E_{\text{xc}}[\rho]}{\delta \rho(\mathbf{r})}(\mathbf{r}) \quad (66)$$

↑ Hartree potential ↑ Exchange-correlation potential

How to solve? See Hartree-Fock (later)

Recent developments in LIBXC — A comprehensive library of functionals for density functional theory



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ABSTRACT

LIBXC is a library of exchange–correlation functionals for density-functional theory. We are concerned with semi-local functionals (or the semi-local part of hybrid functionals), namely local-density approximations, generalized-gradient approximations, and meta-generalized-gradient approximations. Currently we include around 400 functionals for the exchange, correlation, and the kinetic energy, spanning more than 50 years of research. Moreover, LIBXC is by now used by more than 20 codes, not only from the atomic, molecular, and solid-state physics, but also from the quantum chemistry communities.

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Idea

- Assumes (almost) same density everywhere
- Exact for the uniform electron gas *Jellium*
- Used for materials, deprecated
- Strictly local

$$E_{xc}^{\text{LDA}} = \int d\mathbf{r} \rho(\mathbf{r}) \epsilon_{xc}^{\text{LDA}}[\rho(\mathbf{r})] \quad (67)$$

Popular choice: VWN

Idea

- Includes local gradient
- *semi-local* nature: gradient tells a bit about close proximity
- Widely used, both for materials and molecules

$$E_{xc}^{\text{GGA}} = \int d\mathbf{r} \rho(\mathbf{r}) \epsilon_{xc}^{\text{GGA}}(\rho(\mathbf{r}), \nabla \rho(\mathbf{r})) \quad (68)$$

Popular choice: PBE

Idea

- Next derivative and/or kinetic energy density τ
- Sensitive to numerical integration through Laplacian
- Default settings sometimes not good enough

$$E_{xc}^{\text{GGA}} = \int d\mathbf{r} \rho(\mathbf{r}) \epsilon_{xc}^{\text{GGA}}(\rho(\mathbf{r}), \nabla\rho(\mathbf{r}), \nabla^2\rho(\mathbf{r}), \tau), \quad \tau \equiv \sum_i^{\text{occ}} |\nabla\psi_i(\mathbf{r})|^2 \quad (69)$$

Popular choices: SCAN (built on constraints!), M06, TPSS

Idea

- Commonly improved band gaps and relative energies between spin states
- Popular choices: B3LYP, PBE0

$$E_{xc}^{\text{hyb}} = a E_x^{\text{HF}} + (1 - a) E_x^{\text{DFT}} + E_c^{\text{DFT}} \quad (70)$$

$$E_x^{\text{HF}} = -\frac{1}{2} \sum_{i,j} \iint \psi_i^*(\mathbf{r}_1) \psi_j^*(\mathbf{r}_2) \frac{1}{r_{12}} \psi_j(\mathbf{r}_1) \psi_i(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2 \quad (71)$$

Idea

- Computationally more efficient
- Popular choice: HSE06

$$\frac{1}{r_{12}} = \frac{1 - [\alpha + \beta(1 - \omega_{\text{RSF}}(\gamma, r_{12}))]}{r_{12}} + \frac{\alpha + \beta(1 - \omega_{\text{RSF}}(\gamma, r_{12}))}{r_{12}} \quad (72)$$

Can be tuned in many cases to specific system based on Koopman's condition

$$-\varepsilon_{\text{HOMO}}(N_e) = \text{IE}(N_e) = E(N_e - 1) - E(N_e) \quad (73)$$

Implications

- Replace the $3N$ -dimensional wave function with 3D electron density.
- There is one functional (function of functions) yielding the energy, hence the name density functional theory: $E[\rho]$
- This functional is...
 - Universal: Does not depend on the system!
 - Unknown: Topic of current research
- We can obtain the electron density by minimizing the energy.