Born-Oppenheimer Approximation

Formal picture 11

External

Nuclear repulsion Electron-nuclei attraction Electronic repulsion

Formal picture 12

$\hat{\mathrm{H}}(\mathbf{r}_i,\mathbf{R}_I)=\hat{\mathrm{T}}_{\mathrm{n}}(\mathbf{R}_I)+\hat{\mathrm{T}}_{\mathrm{e}}(\mathbf{r}_i)+\hat{\mathrm{V}}_{\mathrm{ee}}(\mathbf{r}_i)+\hat{\mathrm{V}}_{\mathrm{en}}(\mathbf{r}_i,\mathbf{R}_I)+\hat{\mathrm{V}}_{\mathrm{nn}}(\mathbf{R}_I)+\hat{\mathrm{V}}_{\mathrm{ext}}$

What if Hamiltonian would be separable into electronic and nuclear parts?

Equipartition theorem (thermodynamics):

The time-averaged kinetic energy in a degree of freedom only depends on the temperature.

$$
E_{\text{kin,n}} \approx E_{\text{kin,e}}
$$

\n
$$
\Rightarrow m_n v_n^2 \approx m_e v_e^2
$$

\n
$$
\Rightarrow \frac{v_e}{v_n} \approx \sqrt{\frac{m_n}{m_e}} \qquad \frac{m_n}{m_e} \approx 1800
$$

Separation: slow nuclei and fast electrons live on different time scales.

Born-Oppenheimer approximation and the state of the state o

Approximation: Hamiltonians and wave functions of the two are separable by time scales.

$$
\hat{\mathrm{H}}(\mathbf{r}_i, \mathbf{R}_I) \equiv \hat{\mathrm{H}}_{\mathrm{e}}(\mathbf{r}_i | \mathbf{R}_I) + \hat{\mathrm{H}}_{\mathrm{n}}(\mathbf{R}_I)
$$

$$
\Psi(\mathbf{r}_i,\mathbf{R}_I)\equiv\Psi_{\mathrm{e}}(\mathbf{r}_i|\mathbf{R}_I)\Psi_{\mathrm{n}}(\mathbf{R}_I)
$$

$$
f(x|b) = f_b(x) \equiv \log_b x
$$

Parameter (static)

Argument (changeable)

- 1. Consider nuclei fixed (and potentially classically).
- 2. Solve Schrödinger equation for electrons only.

$$
\hat{\mathrm{H}}_{\mathrm{e}}(\mathbf{r}_i|\mathbf{R}_I)\Psi(\mathbf{r}_i|\mathbf{R}_I)=E\Psi(\mathbf{r}_i|\mathbf{R}_I)
$$

System defines molecular Hamiltonian \hat{H}

 $\hat{H} = \hat{T} + \hat{V} = \hat{T}_{e} + \hat{V}_{en} + \hat{V}_{ee} + \hat{V}_{nn} + \hat{V}_{ext}$

$$
-\sum_{i,j}\frac{Z_ie^2}{4\pi\varepsilon_0\left|\mathbf{R}_i-\mathbf{r}_j\right|}
$$

Coulomb interaction electrons-nuclei

$$
+\sum_{i,j>i}\frac{e^2}{4\pi\varepsilon_0\left|\mathbf{r}_i-\mathbf{r}_j\right|}
$$

Coulomb interaction electrons-electrons

$$
+ \sum_{i,j>i} \frac{Z_i Z_j e^2}{4\pi\varepsilon_0 |\mathbf{R}_i - \mathbf{R}_j|}
$$

Coulomb interaction nuclei-nuclei

Wave function

$$
\Psi=\Psi({\bf r}_1,{\bf r}_2,{\bf r}_3,\ldots,{\bf r}_n)
$$

Wave function and the set of the se

$$
\Psi=\Psi({\bf r}_1,{\bf r}_2,{\bf r}_3,\ldots,{\bf r}_n)
$$

Impractical:

- 3n dimensions for n electrons
- Almost always no closed form expression
- Numerical solution: discretisation in basis functions

$$
\text{Costly:} \quad [\text{s}] \simeq \exp(2n)/10^4 \qquad [\text{MB}] \simeq \exp(n)/2
$$

Methane $CH₄$

https://gist.github.com/ferchault/1e6009e78310200673969ef96c9dad1d

- 1. In principle, properties could be done arbitrarily accurately, but that is too expensive.
- 2. Two core approximations:

Electrons feel static nuclei and nuclei feel averaged electrons. ("Born-Oppenheimer")

Finite accuracy of wave function / numerics. ("Discretisation")

3. Mental picture: static atoms and electron density.